ESRF	Experiment title: FLOW PROPERTIES OF SOFT CRYSTALS : THE CASE OF A BODY CENTERED EUDIC PHASE	Experiment number: SC 401
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Report: fig 1 a) 1 c) 16) fig 2 2000 Stress [Pa] 1000 F68 at T = 27°C 400 fig 3 10⁻³ 10-2 10³ 10-1 10⁰ 10² 10 Shear Rate [s⁻¹]

The goal of the experiment was to compare the behaviour under steady shear of well ordered cubic mesophases of different symmetries: **fcc** and **bcc** respectively. These cubic phases are obtained from weakly amphiphilic triblock copolymers (polyoxyethylene / polyoxypropylene / polyoxyethylene) which spontaneously self assemble in water into spherical micelle having a dense POP core and a diffuse POE corona. At sufficient concentration these micelles pile up into well ordered cubic mesophases. The nature of the lattice depends on the relative length of the POE and POP moieties. In a former experiment conducted in April 1996 we characterized in details the structural evolution under shear of the **fcc** cubic phase obtained from the Pluronic F 108-copolymer (see the former experiment report: I lost the number): a very progressive alignment of the scattering pattern could be interpreted in terms of a progressive transformation of the initial polycrystalline texture into the so called "layer sliding" regime.

During run no SC401, we characterized in a similar manner the response to shear of the bcc cubic phase from the Pluronic F68 copolymer. In the radial geometry (X-beam along the direction of the velocity gradient of the shear cell), we distinguish two distinct regimes depending on the imposed shear rate; below 1 s^{-1} , we observe a weakly oriented powdered texture for which the bragg spots have a banana like angular distribution (fig (1a): above 100 s⁻¹, we have a fully oriented sample (fig. [c) with Bragg peaks having sharp angular positions: this is the layer sliding regime for the **bcc** structure. In the cross over domain $(1 \text{ s}^{-1} \text{ to } 100 \text{ s}^{-1})$ between these two pure regimes, the scattering patterns (fig 1b) consist of the weighted sum of the signatures of the two pure regimes: so clearly we have coexistence of the powdered texture and the layer sliding texture in distinct bands inside the sheared sample. Interestingly, the texture coexistence precisely correspond to the shear rate range where the stress response to increasing shear rates levels off and stays nearly constant (fig 2). This texture coexistence at constant stress is thus very analogous to the phase coexistence at constant chemical potential in a binary mixture undergoing a first order phase separation. In this respect the behavior of the bcc structure differs markedly from that of the fcc structure for which the structural change upon increasing rates is very progressive. This very interesting analogy with a phase transition is under current theoretical investigation.

On the other hand, the layer sliding regime in the present bcc structure appears more complex than expected: the scattering pattern in fig 3 is taken in the tangential geometry (beam along the velocity direction) at high shear rate. It complexity implies that two non symmetrically equivalent orientations are simultaneously present in the high shear regime. The interpretation of this complex pattern is under progress.