ESRF	<b>Experiment title:</b> In-situ study of aerogelation in a flame soot aerosol	Experiment number: SC 1961
Beamline:	Date of experiment:	Date of report:
ID2	from: 15.06.2006 to: 19.06.2006	07.02.2007
Shifts:	Local contact(s):	Received at ESRF:
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## **Report:**

The aim of the present experiment was to investigate aerogelation of soot particles in an acetylene flame by (ultra) small-angle X-ray scattering [1]. As compared to previous studies, we have attempted to record the scattered intensity over a wider range of scattering wave vector (q),  $10^{-3}$  nm<sup>-1</sup>  $\leq q \leq 0.5$  nm<sup>-1</sup>, by combining earlier experiments recorded at the ID2 pinhole SAXS instrument with Bonse-Hart USAXS measurements. The significantly reduced instrumental background of the Bonse-Hart setup using specially fabricated analyzer crystals facilitated this experiment [2]. Quantitative data analysis was performed using the unified scattering function involving up to three structural levels for primary particles, aggregates and agglomerates [3]. Each structural level provides corresponding radius of gyration ( $R_g$ ), polydispersity ( $\sigma$ ), number density (N), and fractal dimension  $(d_t)$ .

The diffusion flame was produced by a home-made stainless steel cylindrical burner fueled with acetylene. To stabilize the flame, a concentric flow of nitrogen gas was



**Figure 1.** Comparison of Bonse-Hart USAXS and pinhole SAXS data for soot particles in a diffusion flame (acetylene, 130ml/min,  $\varphi \sim 10^{-6}$ ) at different heights above the burner. The extended low *q* range of the USAXS data is essential for characterizing the structure of aggregates.

maintained around the acetylene stream by an outer concentric tube. The spatial stability of the flame was critical for the USAXS measurements which had a typical scanning time of 4 minutes. The flow rate of acetylene was varied between 90 ml/min and 340 ml/min, extending into the regime of aerogelation of soot [4]. The burner was mounted on a sample stage with motorized horizontal and vertical translations. The particle growth in the flame as a function of height above the burner (*HAB*) was probed by moving the burner in the vertical direction. The temperature profile of the flame was recorded using a miniature thermocouple. Using the measured temperature profile and assuming ideal gas behavior, *HAB* was converted to the residence time or kinetic time ( $t_R$ ).

The size  $(R_{gl})$  and polydispersity of compact primary particles and their fractal aggregates as a function of *HAB* was probed using the pinhole SAXS instrument. It is rather puzzling that the growth of primary particles follows an Avrami type beavior typical of classical nucleation and growth [5] resulting in a compact

morphology with terminal radius of gyration of  $R_{gl}(\infty) \approx 27$  nm. In addition the growth exponent, *n*, was found to be 1, implying a linear growth of primary particles. This is suggestive of a scenario where primary particles are formed by the stacking of lower aspect ratio polycyclic aromatic hydrocarbons [6].

However, the limited low q range of the SAXS data recorded at a burner to detector distance of 10 m did not allow to fully characterize the aggregates and their final size. Therefore, SAXS results can only represent a lower limit of the aggregate  $R_{g2}$ .

During the present experiment, further details about the aggregation process of the primary particles could be monitored in the USAXS data which extends to the low q limit of about 0.001 nm<sup>-1</sup>. Figures 1 shows the typical evolution of I(q) at different *HAB* for a flow rate of 130 ml/min. For a comparison, corresponding SAXS data are also plotted. At lower *HABs*, only the scattering from the primary particles and the aggregates is observable. The low q upturn for *HAB*  $\geq$  14 mm indicates the presence of agglomerates which can be described only by including a third structural level. The terminal  $R_{g2}$  of aggregates levels off at about 250 nm with  $d_f \sim 1.96$ . Similar features observed at a lower flow rate of 90 ml/min as a function of *HAB* and residence time  $t_R$  are shown in Figure 2.

Figure 3 presents the evolution of  $R_{g2}$  of the aggregates obtained from the SAXS and USAXS data. Within the error bars, the time dependence is consistent with  $t_R^{1/2}$  growth law. The corresponding fractal dimension  $d_f \sim 2$  without any systematic dependence on  $t_R$  signifies a constant mechanism of aggregation through the flame. In diffusion limited clustercluster aggregation,  $R_g \sim t^{z/d_f}$ , with z=1. Therefore,  $t_R^{1/2}$ power-law growth presented in Fig. 3 is an indicative of the diffusion limited aggregation process.

Although the  $R_{g3}$  of agglomerates cannot be completely estimated from the USAXS data, the increased upturn at low q (see Figs. 2 and 3) indicates that the corresponding  $d_f$  is in the range of 2 to 2.4. This large value of  $d_f$  is attributed to shear flow effects and percolation of clusters. The size of the agglomerates is at least an order of magnitude larger than the aggregates but with a more compact morphology. The obtained structural parameters display similar dependence on the residence time for different flow rates investigated suggesting a constant growth mechanism.



**Figure 2.** Typical evolution of scattered intensity as a function of *HAB* and residence time  $t_R$  for an acetylene flow rate of 90 ml/min. The overall features are similar to that shown in Fig. 2.



**Figure 3.** Radius of gyration of fractal aggregates as a function of the residence time in the flame for three different flow rates.  $R_{g2}$  has been deduced from independent fits to SAXS and USAXS data and the terminal value is more reliably deduced from USAXS. The solid line corresponds to  $t^{1/2}$  behavior which is consistent with a diffusion limited growth.

Aerogelation was found to be a macroscopic phenomenon without significant changes in the aggregate or agglomerate level.

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