



**Experiment title:**  
Imaging of the energy flow in the three-beam case of interference

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HS 722

**Beamline:**  
ID 22

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10

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**Report:**

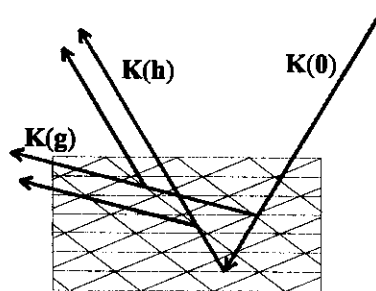


Fig. 1: three-beam case in the direct space (crystal)

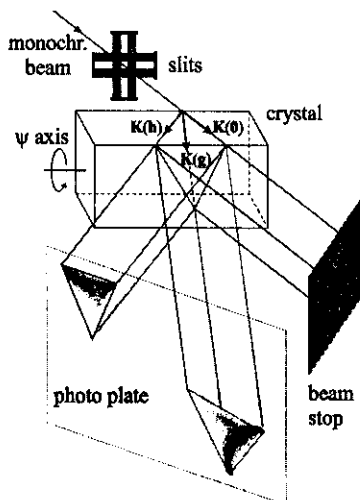


Fig. 2: experimental setup

The purpose of this experiment was the systematical investigation of the energy flow in the three-beam case of x-ray diffraction. In this case, two Bragg reflections  $\mathbf{h}$  and  $\mathbf{g}$  are simultaneously excited and three strong wave fields  $\mathbf{K}(\mathbf{0})$ ,  $\mathbf{K}(\mathbf{h})$ ,  $\mathbf{K}(\mathbf{g})$  exist inside the crystal. The intensity diffracted into the direction of  $\mathbf{K}(\mathbf{h})$  can then be depicted as an interference between the wave diffracted at the lattice planes of  $\mathbf{h}$  (direct wave) and a wave which is diffracted at  $\mathbf{g}$  and afterwards at  $\mathbf{h}-\mathbf{g}$  (Umweg wave). /3/ For the systematical investigation of such a case, a  $\Psi$ -rotation (a rotation around the primary diffraction vector  $\mathbf{h}$ ) was used. Thus, the distance of the reciprocal lattice point  $\mathbf{G}$  to the Ewald sphere can be changed independently of  $H$ .

In the two-beam case the energy flow is limited inside the Borrmann triangle formed by the direction of  $\mathbf{K}(\mathbf{0})$  and  $\mathbf{K}(\mathbf{h})$ . For three strong waves in a plane-parallel crystal plate the corresponding figure is a pyramid. Its top is given by the incidence point of radiation on the crystal and a triangular base on the exit surface. The sides of the pyramid are limited by  $\mathbf{K}(\mathbf{0})$ ,  $\mathbf{K}(\mathbf{h})$  and  $\mathbf{K}(\mathbf{g})$ , respectively. All directions inside this pyramid which start from the top are possible directions of energy flow. For the imaging of the intensity distribution in the triangular basis we used a kind of pinhole topography with a cross section of the incidence beam reduced by slits in both dimensions to a 'point' of  $30 \times 30 \mu\text{m}^2$  (see Fig. 2).

With this experimental setup we could take several systematical series of topographs in different three-beam interferences. The  $\psi$  - position of the crystal was the control parameter. Because of the precise and fast six-circle diffractometer at ID22 we could reach a better relation between the topograph and the  $\psi$  - position compared with our first attempt at the end of HS-401. /1/

The results of our experiment show that the topographs image the dispersion surface. This is confirmed by the investigation of a three-beam case with a forbidden coupling reflection  $h-g$ . Although there is no direct energy transfer between the  $h$  and  $g$  reflection we found an intensity spread over the complete triangular basis of the pyramid in the exact three-beam position (see Fig. 3 left). The contrast change in our topographic series in dependence on  $\psi$  describes clearly the intensity transfer between the reflections  $h$  and  $g$  when  $G$  passes the Ewald sphere. A special characteristic feature in all of our pinhole topographs is the visibility of the inflection point on the dispersion surface (see Fig. 3 right). As another important point in our results we could observe a systematic in the relation between the shape of the diffraction pattern and the triplet-phase of the reflections.

In a last step before the end of our beam time we had the opportunity to change the experimental setup into a  $90^\circ$  reflection position. In spite of the polarisation obliteration we observed intensity.

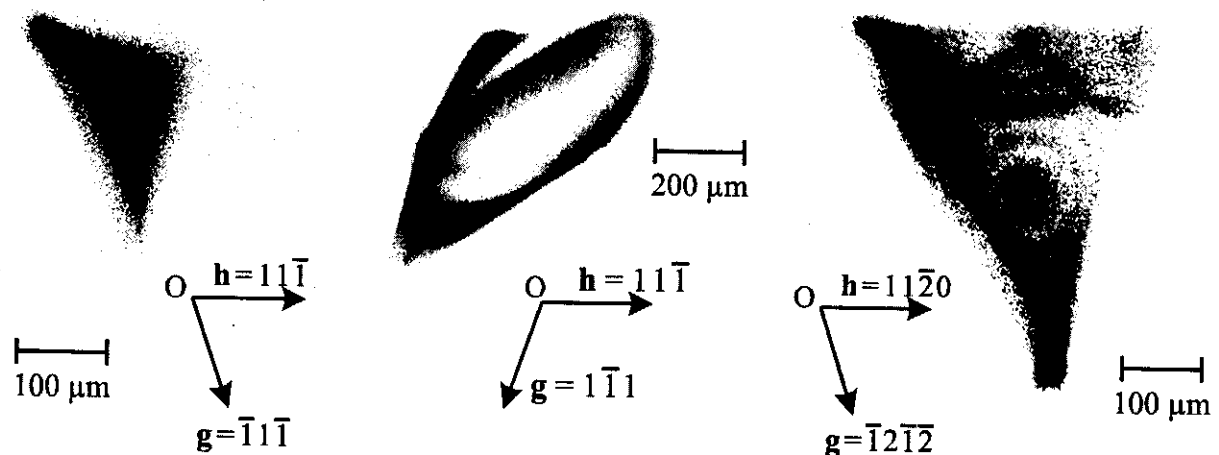


Fig. 3: Examples for pinhole topographs which show the change in the energy distribution for different three-beam cases. The excited three-beam cases are put as inscriptions in the figures. Each of the pictures shows the respective  $g$  reflection.

left: A 1 mm thick silicon crystal plate with a  $(111)$  surface is shown. In this case the structure factor of the coupling reflection  $h-g = (200)$  is zero.

center: This picture presents a 2 mm thick silicon crystal plate with a  $(111)$  surface taken on nuclear plates with a strong coupling reflection  $h-g = (02\bar{2})$ . The triplet phase in the centrosymmetric silicon is zero.

right: A 1 mm thick quartz crystal with a  $(\bar{1}100)$  surface and an incident beam with the wavelength of  $\lambda \approx 0.749 \text{ \AA}$  was used. Here the triplet phase has a value of about  $30^\circ$ .

#### references

- /1/ F. Heyroth, H.-R. Höche and C. Eisenschmidt „Imaging of the energy flow in the three-beam case of X-ray diffraction“ J. Phys. D: Appl. Phys. 32 (1999), A133 - A138
- /2/ F. Heyroth, H.-R. Höche and C. Eisenschmidt „Contrast in X-ray section topographs of perfect silicon crystals using the Laue-Laue three-beam case of diffraction“ J. Appl. Cryst. 32 (1999), 489 - 496
- /3/ Weckert, E. & Hümmer, K. (1997) „Multiple-Beam X-ray Diffraction for Physical Determination of Reflection Phases and its Applications“ Acta Cryst. A53, 108-143.