



**Experiment title:**  $^{57}\text{Fe}$  NIS studies of phonon DOS in superconducting FeSe systems as function of pressure and composition

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HS-4098

**Beamline:**  
ID18

**Date of experiment:**  
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Shifts: 18

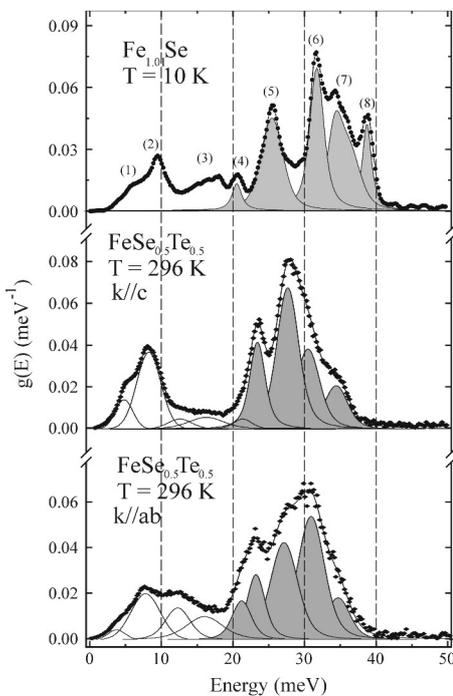
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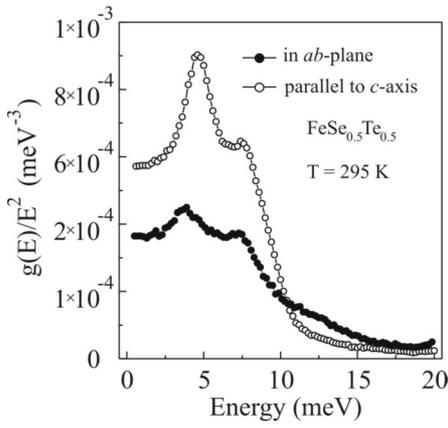
**Report:**



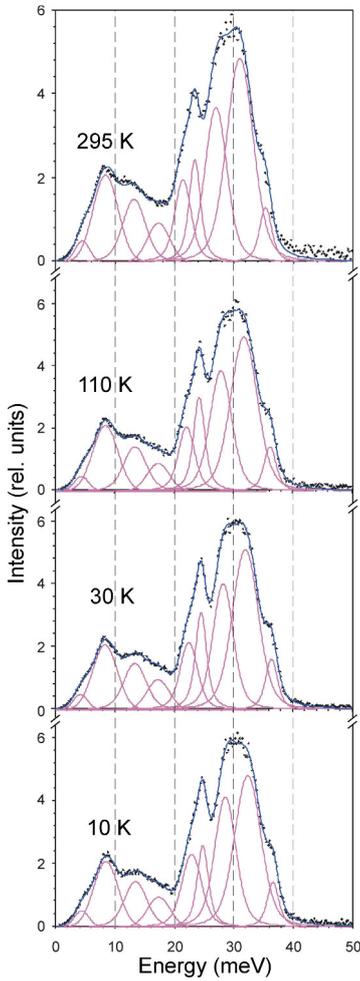
**Fig.1:** Fe partial phonon-DOS for  $\text{Fe}_{1.01}\text{Se}$  at 10 K ([7]) and of  $\text{FeSe}_{0.5}\text{Te}_{0.5}$  X-tals at RT measured with different orientation of the SR k-vector.

$\text{FeSe}$  and  $\text{FeSe}_{0.5}\text{Te}_{0.5}$  are model systems for the new high- $T_C$  superconductors with  $\text{FeX}$  ( $X = \text{As}$  and  $\text{Se}$ ) layers. In the previous beamtime HS-3978, we studied in detail the temperature and pressure dependence of the local Fe phonon-DOS in a superconducting (s.c.)  $\text{Fe}_{1.01}\text{Se}$  sample, well characterized and intensively investigated by the applicants in cooperation with Princeton University and Max-Planck Institute of Chemistry, Mainz [1-5].  $\text{FeSe}$  could be, because its much simpler structure but still with the same  $\text{FeX}$  layers a clue compound [6] for the understanding of the principal mechanisms of superconductivity in these  $\text{FeX}$ -based systems. In particular, the dramatic increase of  $T_C$  from 8 K at ambient pressure to 37 K at 9.0 GPa [2] points to a pressure modulated mechanism for superconductivity in  $\text{FeSe}$  systems. Similar to the s.c.  $\text{FeAs}$  systems  $\text{Ba}_{1-x}\text{K}_x\text{Fe}_2\text{As}_2$  and  $\text{SmFeAsO}_{1-x}\text{F}_x$ , a large isotope effect on  $T_C$  has been detected for the  $\text{FeSe}$  system [6]. Therefore a detailed study of the local phonon density-of-states (DOS) at the Fe sites, which are here the active “players”, is of actual scientific interest [6]. We have just published the results of the  $^{57}\text{Fe}$ -NIS study on the temperature and pressure dependence of the partial phonon-DOS in  $\text{Fe}_{1.01}\text{Se}$  [7].

In the present beamtime we studied in detail the phonon-DOS in  $\text{FeSe}_{0.5}\text{Te}_{0.5}$  as function of temperature and, by the use of single crystals, the polarisation dependence of the phonon modes. We applied again  $^{57}\text{Fe}$  nuclear inelastic scattering (NIS) spectroscopy, with the high energy resolution of 0.7 meV for the 14.413 keV radiation, provided at beamline ID18. The present sample, 100% enriched in  $^{57}\text{Fe}$ , was prepared by Vladimir Tsurkan from the group of Prof. Alois Loidl (Univ. Augsburg), specialised in this system [8]. The investigated sample had a s.c. temperature of  $T_C = 14$  K at ambient pressure. Application of pressure leads to an increase of  $T_C$  up to 25 K at 3 GPa, as determined in Mainz, similar to the reports of other groups [9]. In Fig. 1 we present the phonon-DOS of the  $\text{FeSe}_{0.5}\text{Te}_{0.5}$  X-tal, measured at 296 K with the k-vector of the monochromatized synchrotron radiation parallel and perpendicular to the c-axis of the tetragonal structure. For comparison, we show on top of Fig.1 the phonon-DOS of  $\text{Fe}_{1.01}\text{Se}$ , measured at 10 K in the previous beamtime [7]. We refer to our recent publication for details of the analysis of the pronounced structures of the DOS, labelled 1-3 for acoustic and 4-8 for optical phonon modes [7]. The same principal features are present in the phonon-DOS of the  $\text{FeSe}_{0.5}\text{Te}_{0.5}$  X-tal sample, exhibiting now a pronounced polarization dependence, which allows us to prove (or disprove) our assignment of certain optical modes in our publication, based in part on comparison with theoretical



**Fig.2:** Reduced form of Fe phonon-DOS for  $\text{FeSe}_{0.5}\text{Te}_{0.5}$  X-tals with different orientations.



**Fig.3:** Phonon-DOS of polycryst.  $\text{FeSe}_{0.5}\text{Te}_{0.5}$  at different temperatures.

calculations of the phonon-dispersions and phonon-DOS of FeSe (see refs. in [7]). In addition a pronounced shift to lower energy is observed, different for different modes in the phonon-DOS of  $\text{FeSe}_{0.5}\text{Te}_{0.5}$ , by  $\sim 4\%$  for mode 5, assigned to the  $B_{1g}$  Raman mode, and up to  $13\%$  for the lowest acoustic mode 1, in comparison to  $\text{Fe}_{1.01}\text{Se}$  [7]. These shifts result from the combined effects of a lattice expansion by  $\sim 10\%$  and an increase of the effective mass  $m^*$  by  $\sim 8\%$ , both caused by the substitution of Se by Te. Of course, the highly anisotropic elastic properties of the layered PbO-structure with a  $\sim 3$ -fold stronger compression (expansion) of the  $c$ -axis than the  $a,b$ -plane result in different pressure/volume induced variations of the phonon-mode energies (see below). In addition, the random substitution of the Se sites by Te caused obviously a broadening of certain structures resulting from a combined motion of the Fe and Se,Te atoms, most pronounced in the optical modes 6–8, which also exhibit a lowering of their energies by  $\sim 8\%$ . In contrast, the mode 5, assigned by us to the  $B_{1g}$  mode [7], which originates mainly from Fe atoms vibrating in opposite

direction parallel to the  $c$ -axis, exhibit only a small shift of  $\sim 4\%$  from 24.4 meV (FeSe) to 23.4 meV ( $\text{FeSe}_{0.5}\text{Te}_{0.5}$ ) with no resolvable spectral broadening. This is easily explained by the fact that mainly the Fe atoms are involved in this mode. The assignment to  $B_{1g}$  is further supported by the observation of this mode at 23.0 meV in the Raman spectrum of isostructural  $\text{FeSe}_{0.3}\text{Te}_{0.7}$  [10]. In Fig. 2 we present a plot of the phonon-DOS  $g(E)$  in the reduced form,  $g(E)/E^2$ , exemplifying the strong anisotropy of the acoustic modes, with a “soft” acoustic mode at 4.7(2) meV clearly connected with the soft  $c$ -axis (for details, see [7]). These anisotropic properties are reflected in the derived parameters like the Debye temperatures, the sound velocities and the Lamb-Mössbauer  $f$ -factors, similar to the case of  $\text{Fe}_{1.01}\text{Se}$  [7].

In Fig. 3 we present the phonon-DOS of a polycrystalline, but textured (resembling the  $k//ab$  phonon-DOS)  $\text{FeSe}_{0.5}\text{Te}_{0.5}$  sample at different temperature, e.g. 10 K and 30 K, below and above the s.c. transition, also at 110 K, above the tetragonal-orthorhombic transition and at 295 K for comparison with the oriented X-tal samples shown in Fig. 1. Again we see with decreasing temperature a sharpening of the spectral features, not as pronounced as for  $\text{Fe}_{1.01}\text{Se}$ , and shift of the phonon-modes by  $\sim 4\%$  to higher energies, very similar to  $\text{Fe}_{1.01}\text{Se}$  [7]. Here the data analysis is, as for the data of  $\text{FeSe}_{0.5}\text{Te}_{0.5}$  in Fig. 1, in a preliminary stage. We want to use the information obtainable from difference spectra, evaluated similar as in a  $^{119}\text{Sn}$ -NIS study of SnO with the same PbO structure [11], to search for subtle details in the phonon-DOS above/below  $T_C$  and the tetragonal/orthorhombic transition, respectively. As a preliminary resume, we obtained new and highly informative data from the polarisation-dependent phonon-DOS of X-tal samples of  $\text{FeSe}_{0.5}\text{Te}_{0.5}$  as well as from their temp.-dependent behaviour important for their (possible) connections to the s.c. properties. These new data on  $\text{FeSe}_{0.5}\text{Te}_{0.5}$  will be published together with the Augsburg group, where the well-defined  $^{57}\text{Fe}$ -enriched samples were prepared [12].

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