



	<b>Experiment title:</b> Atomic diffusion in B2-AuMg measured by XPCS	<b>Experiment number:</b> HS-4357
<b>Beamline:</b> ID10	<b>Date of experiment:</b> from: 16.11.2011 to: 22.11.2011	<b>Date of report:</b> 01.09.2016  <i>Received at ESRF:</i>
<b>Shifts:</b> 18	<b>Local contact(s):</b> Dr. Federico Zontone	
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We proposed to study atomic diffusion in  $Au_xMg_{1-x}$  intermetallic alloys by atomic-scale X-ray photon correlation spectroscopy (aXPCS) in order to determine the diffusion mechanisms on both sides of the stoichiometric composition. Measurements on the prepared  $Au_{64}Mg_{36}$  sample were not successful due to an overall low intensity and a very small speckles contrast  $\beta$  of the scattered X-rays. This was because the sample, which was prepared as thin as technically possible considering the softness of this single crystal, was a very strong absorber of 8 keV radiation. Instead, we measured during the beamtime HC-885 in April 2013 an equivalent B2-ordered phase  $Ag_{58}Mg_{42}$ . The speckles contrast  $\beta$  was significantly higher due to some improvements in the beamline setup.

The single crystal of  $Ag_{58}Mg_{42}$  intermetallic alloy was oriented with its [110] direction normal to the surface. All measurements were done using the iKon-M CCD chip with a sample-to-detector distance of 67 cm. Exposure times were between 4 s for small angles and 8 s for angles larger than  $2\theta = 8^\circ$ . The illuminated area on the sample was determined by slits of  $20 \times 20 \mu\text{m}^2$ . Measurements were taken at 423 K.

B2 alloys are usually considered either to be of the triple-defect type, where antisite atoms exist only on one sublattice and vacancies only on the other, or to be of the anti-structure type, where constitutional defects are formed as antisites for an excess A or B atoms, respectively [1]. Ag-

Mg alloys are usually assigned to the second class [2] but some authors report a hybrid behavior. The best fit was achieved mixing  $\langle 1\ 1\ 0 \rangle$  and  $\langle 1\ 0\ 0 \rangle$  effective jumps in the proportion 3:1. Our measurements suggest that chemical short-range order is rather weak in this system, thus we assumed that the SRO is constant  $I_{\text{SRO}}(\mathbf{q}) = 1$ . The resulting simultaneous fit to all data, i.e. to the azimuthal and to the polar scans is shown in Fig. 1.

As can be seen from Fig. 1,  $\langle 1\ 1\ 0 \rangle$  jumps are dominant in  $\text{Ag}_{58}\text{Mg}_{42}$  with and addition of  $\langle 1\ 0\ 0 \rangle$  jumps. The diffusion coefficient at 423 K is  $2.54(9) \times 10^{-23} \text{ m}^2\text{s}^{-1}$  [3].

Concluding, the microscopic diffusion mechanism in Ag–Mg is similar to that of Fe–Al alloy at high temperatures (above 1200 K) [4] but definitively different than in Fe–Al at about 700 K [5]. One can presume that similar atomistic mechanism is responsible for diffusion in Ag–Mg and in moderately ordered Fe–Al [4]. This mechanism takes place via correlated exchanges of a vacancy leading to farther effective jumps of the majority atoms. The precise ratio of frequencies of  $\langle 1\ 1\ 0 \rangle$  to  $\langle 1\ 0\ 0 \rangle$  jumps depends on the particular values of the exchange energy between both atomic species and a vacancy. This is definitely different than in Fe–Al [4], where  $\langle 1\ 1\ 0 \rangle : \langle 1\ 0\ 0 \rangle$  was 2:1 instead of 3:1 in an Ag–Mg intermetallic alloy.

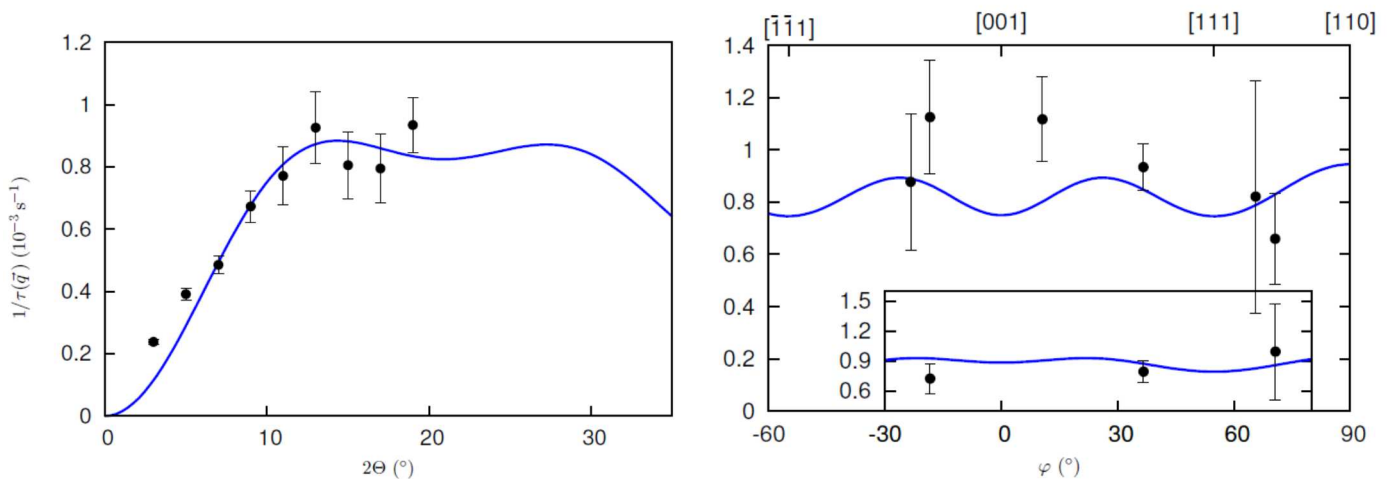


Fig. 1 (left) Inverse correlation time as a function of scattering angle  $2\theta$ ; (right) as a function of azimuthal angle  $\varphi$  for  $2\theta = 19^\circ$  and  $2\theta = 17^\circ$  (insert). The simultaneous fit with an atomistic jump model described in the text is shown by the blue line. Figures taken from M. Stana PhD Thesis [3].

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## References

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